

EQUATION OF STATE FOR A BOILING LIQUID

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Universal dimensionless-pressure functions are found which can be used to estimate the error involved in approximating the equation of state as a linear dependence of the density of the liquid-vapor mixture on its enthalpy.

In studies of processes in steam generators, the equation of state of the boiling liquid is sometimes represented as a linear dependence of the specific volume  $v$  on the enthalpy [1, 2]:

$$v = a + bi, \tag{1}$$

where the coefficients  $a$  and  $b$  are assumed to depend weakly on the pressure (or the boiling temperature). Let us determine the form of these dependences; if  $v = v' + x(v'' - v')$  and  $i = i' + xr$ , we have

$$a = v' - \frac{v'' - v'}{r} i', \tag{2}$$

$$b = \frac{v'' - v'}{r}. \tag{3}$$

It is sufficiently accurate to assume that the physical parameters of the liquid and vapor on the saturation curve, i.e., the liquid enthalpy  $i$ , the specific volumes  $v'$  and  $v''$  of the liquid and vapor, and the heat of vaporization  $r$ , are single-valued functions of only the pressure [3] or only the temperature [4]. The value of any parameter can be represented as the product of some constant for the given liquid and parameter and a function (universal for a group of thermodynamically similar substances) of the dimensionless temperature  $\bar{T} = T/T_{cr}$  or the dimensionless pressure  $\bar{p} = p/p_{cr}$  ( $T_{cr}$  and  $p_{cr}$  are the critical values):

$$\chi = \chi_* \pi_\chi.$$

Graphs of the universal function  $\pi_\chi$  and the values of the constant coefficients  $\chi_*$  for several thermodynamically similar substances can be found in the articles by Borishanskii [3] and Povarin [4]. Approximation of these functions by analytic expressions yields:

$$\begin{aligned} \pi_v &= v_1 + v_2 \bar{T}, \\ \pi_v'' - \pi_v' &= \pi_v''' = v_3 \bar{T}^n, \\ \pi_r &= r_1 - r_2 \bar{T}^2, \\ \pi_i &= i_1 \bar{T} + i_2 \bar{T}^2. \end{aligned}$$

Here we have arbitrarily equated the initial enthalpy to zero; this has no fundamental bearing on this discussion. The constant coefficients in these expressions for a group of substances thermodynamically similar to water (see [3, 4] for a definition of these substances) are:  $v_1 = 0.03968$ ;  $v_2 = 0.08445$ ;  $v_3 = 0.55$ ;  $n = -10.075$ ;  $r_1 = 4.85$ ;  $r_2 = 3.27$ ;  $i_1 = 3$ ; and  $i_2 = 1.85$ .

On the basis of the above discussion, we can now obtain without difficulty the following expressions for  $a$  and  $b$ :

$$a = v_* \left( \pi_v' - \frac{i_*}{r_*} \frac{\pi_v''' \pi_i'}{\pi_r} \right), \tag{4}$$

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$$b = \frac{v_*}{r_*} \frac{\pi_v'''}{\pi_r}. \quad (5)$$

Calculations for water show that the first term in (4) has essentially no effect on  $a$  up to  $\bar{T} = 0.8$ , and that its effect becomes significant only near the critical parameters (at  $\bar{T} > 0.9$ , the first term increases, exceeding 4.4%). We therefore assume that

$$a = - \frac{v_* i_*}{r_*} \frac{\pi_v'''}{\pi_r} \frac{\pi_i'}{\pi_i}.$$

The coefficients  $a_* = v_* i_* / r_*$  and  $b_* = v_* / r_*$  depend only on the nature of the substance, and can be calculated by any working medium. We can therefore find the following functions of the dimensionless parameter  $\bar{p}$  (or  $\bar{T}$ ) from (5) and (6):

$$\begin{aligned} \pi_a &= - \frac{\pi_v'''}{\pi_r} \frac{\pi_i'}{\pi_i}, \\ \pi_b &= \frac{\pi_v'''}{\pi_r}. \end{aligned} \quad (7)$$

The calculated results show that  $\pi_a$  and  $\pi_b$  can be represented as power-law dependences:

$$\begin{aligned} \pi_a &= a_1 \bar{p}^n, \\ \pi_b &= b_1 \bar{p}^m, \end{aligned}$$

where we have  $a_1 = 1.7$ ,  $n = -0.75$ ,  $b_1 = 0.31$ , and  $m = -1$  for the group of substances which are thermodynamically similar to water (at pressures below 30 bars, the error in the  $\pi_b$  determination of these  $b_1$  and  $m_1$  values increases, exceeding 4%). For the group of metals defined in [5], we have  $a_1 = -0.86$ ,  $n = -0.19$ ,  $b_1 = 0.67$ , and  $m = -0.82$ .

The exponents for both functions are seen to be significantly different from zero, so the assumptions that  $n = -1$  and  $m = -1$  give a more accurate result (see [1]).

Accordingly, we can use functions (7) to find the changes in  $\pi_a$  and  $\pi_b$  (which are proportional to  $a$  and  $b$ ) over the pressure range in a steam-generator channel, and finally estimate the error in the determination of the specific volume; for this purpose, it is more convenient to use the equation of state (1) converted to the form

$$\bar{v} = \frac{i_*}{r_*} (\pi_a + \pi_b \bar{i}),$$

where

$$\bar{v} = v/v_*, \quad \bar{i} = i/i_*.$$

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